



# PRECISE MOMENT ASYMPTOTICS FOR THE STOCHASTIC HEAT EQUATION OF A TIME-DERIVATIVE GAUSSIAN NOISE\*

Heyu LI (李贺宇)

*School of Mathematics, Jilin University, Changchun 130012, China*  
*E-mail: heyul15@mails.jlu.edu.cn*

Xia CHEN (陈夏)<sup>†</sup>

*School of Mathematics, Jilin University, Changchun 130012, China;*  
*Department of Mathematics, University of Tennessee, Knoxville TN 37996, USA*  
*E-mail: xchen@math.utk.edu.cn*

**Abstract** This article establishes the precise asymptotics

$$\mathbb{E}u^m(t, x) \quad (t \rightarrow \infty \quad \text{or} \quad m \rightarrow \infty)$$

for the stochastic heat equation

$$\frac{\partial u}{\partial t}(t, x) = \frac{1}{2}\Delta u(t, x) + u(t, x)\frac{\partial W}{\partial t}(t, x)$$

with the time-derivative Gaussian noise  $\frac{\partial W}{\partial t}(t, x)$  that is fractional in time and homogeneous in space.

**Key words** Stochastic heat equation; time-derivative Gaussian noise; Brownian motion; Feynman-Kac representation; Schilder's large deviation

**2010 MR Subject Classification** 60J65; 60K37; 60H15; 60G60; 60F10

## 1 Introduction

Moment asymptotics for solutions to stochastic partial differential equations are known as the problem of intermittency that has been studied extensively in the past two decades [1, 8]. In this work, we investigate the asymptotics problem

$$\mathbb{E}u^m(t, x) \quad (t \rightarrow \infty \quad \text{or} \quad m \rightarrow \infty)$$

for the stochastic heat equation

$$\begin{cases} \frac{\partial u}{\partial t}(t, x) = \frac{1}{2}\Delta u(t, x) + u(t, x)\frac{\partial W}{\partial t}(t, x) & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d \\ u(0, x) = u_0(x) \end{cases} \quad (1.1)$$

\*Received January 27, 2018. Research partially supported by the “1000 Talents Plan” from Jilin University, Jilin Province and Chinese Government, and by the Simons Foundation (244767).

<sup>†</sup>Corresponding author

with the Gaussian noise  $\frac{\partial}{\partial t}W(t, x)$  that is formally given as the time derivative of the mean-zero Gaussian field  $\{W(t, x); (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d\}$  with the covariance function

$$\text{Cov}\left(W(t, x), W(s, y)\right) = \frac{1}{2}(t^{2H_0} + s^{2H_0} - |t - s|^{2H_0})\Gamma(x, y) \quad (t, x), (s, y) \in \mathbb{R}^+ \times \mathbb{R}^d,$$

where the time Hurst parameter  $H_0 \in (0, 1)$  and we assume that the space covariance function  $\Gamma(x, y)$  is locally bounded and has the homogeneity in the sense that

$$\begin{cases} \Gamma(Cx, Cx) = |C|^{2H}\Gamma(x, x) \\ \Gamma(x, x) + \Gamma(y, y) - 2\Gamma(x, y) = \Gamma(x - y, x - y) \end{cases} \quad (1.2)$$

for any  $x, y \in \mathbb{R}^d$  and  $C \in \mathbb{R}$ , where the constant  $H \in (0, 1)$ . Assumption (1.2) can be restated as

$$\begin{cases} \{W(c_0t, cx); (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d\} \stackrel{d}{=} \{c_0^{H_0}c^H W(t, x); (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d\} \\ W(t, x) - W(s, y) \stackrel{d}{=} W(t - s, x - y) \quad (c_0, c > 0, (t, x), (s, y) \in \mathbb{R}^+ \times \mathbb{R}^d). \end{cases} \quad (1.3)$$

For simplicity, we assume that bounded initial condition is as follows:

$$0 < \inf_{x \in \mathbb{R}^d} u_0(x) \leq \sup_{x \in \mathbb{R}^d} u_0(x) < \infty. \quad (1.4)$$

Mathematically,  $\frac{\partial}{\partial t}W(t, x)$  is defined as a generalized centered Gaussian field with

$$\text{Cov}\left(\frac{\partial}{\partial t}W(t, x), \frac{\partial}{\partial s}W(s, y)\right) = \gamma_0(t - s)\Gamma(x, y) \quad (t, x), (s, y) \in \mathbb{R}^+ \times \mathbb{R}^d. \quad (1.5)$$

Here, the time-covariance  $\gamma_0(t - s)$  is morally considered as the derivative

$$\frac{\partial^2}{\partial t \partial s} \left\{ \frac{1}{2}(t^{2H_0} + s^{2H_0} - |t - s|^{2H_0}) \right\}.$$

In particular,

$$\gamma_0(t - s) = \begin{cases} H_0(2H_0 - 1)|t - s|^{-(2-2H_0)} & H_0 > 1/2 \\ \delta_0(t - s) & H_0 = 1/2. \end{cases} \quad (1.6)$$

The function  $\gamma_0(\cdot)$  defined in (1.6) is qualified as covariance function as it is non-negative definite. Indeed, it can be shown that

$$\gamma_0(u) = \frac{\Gamma(2H_0 + 1) \sin(\pi H_0)}{2\pi} \int_{\mathbb{R}} e^{i\lambda u} |\lambda|^{1-2H_0} d\lambda \quad u \in \mathbb{R}. \quad (1.7)$$

When  $H_0 < 1/2$ , the function  $|\cdot|^{-(2-2H_0)}$  is no longer non-negative definite and is not qualified for being a covariance function. As consequence, the covariance function  $\gamma_0(\cdot)$  can not be legally defined by (1.6) when  $H_0 < 1/2$ . As  $H_0 < 1/2$ , the function  $\gamma_0(\cdot)$  is defined as a generalized function given in (1.7). It should be emphasized that  $\gamma_0(\cdot)$  is not defined point-wise when  $H_0 < 1/2$ .

Under suitable conditions (such as the one assumed in our main theorem), the solution to (1.1) yields the following Feynman-Kac formula:

$$u(t, x) = \mathbb{E}_x \left[ u_0(B_t) \exp \left\{ \int_0^t W(ds, B_{t-s}) \right\} \right] \quad (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d, \quad (1.8)$$

where  $\{B_t; t \geq 0\}$  is a  $d$ -dimensional Brownian motion independent of  $\{W(t, x); (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d\}$  with  $B_0 = x$ , and “ $\mathbb{E}_x$ ” stands for the Brownian expectation. We point to [5] and [6] for existing literature.